

ENUMATH

ISCHIA

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**A Wavelet Approximation Method  
for Antenna Problems**

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## Outline

**A simple wavelet-based fast boundary element method  
for time-harmonic Maxwell scattering from open surfaces in  $\mathbb{R}^3$**

- **The Electrical Field Integral Equation:**
  - Quasidiagonalization via Hodge decomposition
  - Reformulations suitable for nodal finite element discretizations
  - The  $(p, \phi, m, \lambda, \alpha)$ -formulation
- **Nodal wavelet discretization:**
  - Construction of a nodal wavelet basis
  - Stability, matrix compression, preconditioning and all that
  - Some problems with open surfaces

# An specimen from electronics

ILLUSTRATION DE L'APPROCHE 2-D

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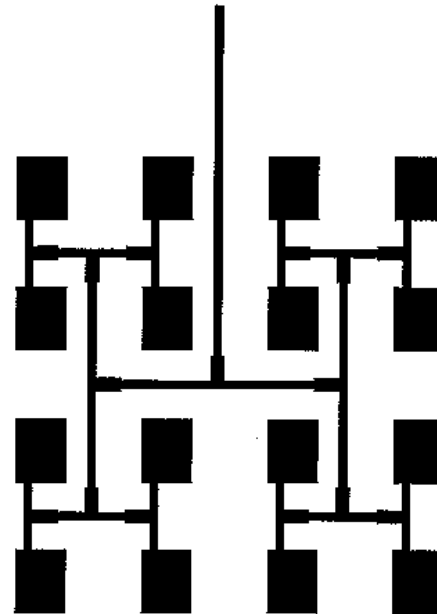


FIG. 2.28: Masque du réseau planaire à 16 éléments.

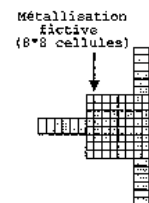


FIG. 2.29: Détail de description de la structure; la ligne quart d'onde.

**A small  
antenna array**

## The geometry

Open surface with p/w smooth boundary  $\Gamma_0 \subset \Gamma$  closed smooth surface  $\subset \mathbb{R}^3$

Simple topologies

Sobolev spaces:  $H^s(\Gamma_0) = H^s(\Gamma)|_{\Gamma_0}$

$$\widetilde{H}^s(\Gamma_0) = \{u \in H^s(\Gamma) \mid \text{supp } u \subset \overline{\Gamma_0}\}$$

Surface differential operators: **grad** =  $\text{grad}_T = -n \times (n \times \text{grad})$

**div** =  $\text{div}_\Gamma = -(\text{grad}_T)^*$

**curl** =  $\text{curl}_\Gamma = -n \times \text{grad}$

**rot** =  $(\text{curl}_\Gamma)^*$

**$\Delta$**  =  $\Delta_\Gamma = \text{grad}_T \text{div}_\Gamma = -\text{rot curl}$

Some spaces on perfectly conducting boundaries:

$$\widetilde{H}_{\text{div}}^{-1/2} = \{u \in \widetilde{H}^{-1/2}(\Gamma_0) \mid \text{div } u \in \widetilde{H}^{-1/2}(\Gamma_0); u \cdot n = 0\}$$

$$H_{\text{curl}}^{-1/2} = \{u \in H^{-1/2}(\Gamma_0) \mid \text{rot } u \in H^{-1/2}(\Gamma_0); u \cdot n = 0\}$$

$$H_{\Delta}^1 = \{u \in H^1(\Gamma_0) \mid \Delta u \in H^{-1/2}(\Gamma_0)\}$$

## The Electrical Field Integral Equation

$$\text{(Maxwell)} \quad \begin{cases} \text{curl } E - i\omega\mu H = 0, & \text{curl } H + i\omega\varepsilon E = 0 \quad \text{in } \mathbb{R}^3 \setminus \overline{\Gamma_0}, \\ E \times n = -E^{\text{in}} \times n & \text{on } \Gamma_0, \quad \text{Silver-Müller r.c.} \end{cases}$$

Single layer potential

$$Vu(x) = \int_{\Gamma_0} \frac{e^{ik|x-y|}}{4\pi|x-y|} u(y) ds(y), \quad k = \omega\sqrt{\varepsilon\mu}$$

Representation formula in  $\mathbb{R}^3 \setminus \Gamma_0$

$$E = \frac{i}{\omega\varepsilon} \text{grad } V \text{ div}_{\Gamma} u + i\omega\mu Vu, \quad u = [H \times n]_{\Gamma_0} \in \widetilde{H}_{\text{div}}^{-1/2}, \quad \text{surface current}$$

Integral equation on  $\Gamma_0$  (EFIE)

$$\boxed{\text{grad } V \text{ div } u + k^2 V_T u = f} = i\omega\varepsilon E_T^{\text{in}}, \quad V_T = \pi_T V = -n \times (n \times V \cdot)$$

## Hodge decomposition and quasi-diagonalization

Problem with **EFIE**: Wrong sign of  $k^2$  !

$V$  : strongly elliptic (pos. def. + compact) pseudodifferential operator of order -1

$$L = \text{grad } V \text{ div} + k^2 V_T \implies (Lu, v) = -(V \text{ div } u, \text{div } v) + k^2 (Vu, v)$$

Hodge decomposition in  $\widetilde{H}_{\text{div}}^{-1/2} = \widetilde{H}^{-1/2}(\text{div}_\Gamma, \Gamma_0) = X^1 \oplus X^2$

$$u = \text{grad } p + \text{curl } \phi, \quad p \in H_\Delta^1, \quad \Delta^{\text{Neu}} p = \text{div } u, \quad \phi \in \widetilde{H}^{1/2}(\Gamma_0)$$

Hodge decomposition in the dual space  $H_{\text{curl}}^{-1/2} = H^{-1/2}(\text{rot}_\Gamma, \Gamma_0)$

$$u' = \text{grad } p' + \text{curl } \phi', \quad \phi' \in \widetilde{H}_\Delta^1, \quad \Delta^{\text{Dir}} \phi' = \text{rot } u', \quad p' \in H^{1/2}(\Gamma_0)$$

$$(u^1, u^2) \quad L \sim \begin{pmatrix} L_{11} & L_{12} \\ L_{21} & L_{22} \end{pmatrix} \quad \begin{array}{l} \text{with } L_{12}, L_{21} \text{ compact} \\ \text{and } L_{22} \text{ and } -L_{11} \text{ strongly elliptic.} \end{array}$$

- All is well for Galerkin methods: Stability, convergence etc...

## (p, φ)-formulation

**Problem for finite element approximations: Discrete Hodge decomposition?**

- Exact H. d.: (p, φ)-formulation
- Approximate H. d.:  $H_{\text{div}}$  elements [Bendali '82, Hiptmair-Schwab '01]
- Approximate H. d.: Approximate or mixed  
(p, φ)-formulation [Safa, Buffa-Co.-Schwab '01]

$$(p, \phi) \quad L \sim \begin{pmatrix} -\Delta V \Delta - k^2 \operatorname{div} V \operatorname{grad} & -k^2 \operatorname{div} V \operatorname{curl} \\ k^2 \operatorname{rot} V \operatorname{grad} & k^2 \operatorname{rot} V \operatorname{curl} \end{pmatrix}$$

**Orders:**  $\begin{pmatrix} 3 & 0 \\ 0 & 1 \end{pmatrix}$  **Problem: Energy space  $H^{3/2} \times \widetilde{H}^{1/2}$ :  $C^1$  elements!**

## Perturbed $(m, \phi)$ -formulation

Introduce the surface charge density  $m = \operatorname{div} u \in \widetilde{H}^{-1/2}$ , so that

$$(p, m) \quad p = (\Delta^{\text{Neu}})^{-1} m \in H_{\Delta}^1 / \mathbb{C}$$

This would lead to a  $(m, \phi)$  formulation containing the operator  $(\Delta^{\text{Neu}})^{-1}$ .

Nothing gained so far.

Let  $\Lambda_J \subset H^1(\Gamma_0)$  be a finite element space and

$p_J \in \Lambda_J$  be defined by the variational solution of (p-m):

$$(p_J, m) \quad p_J = (\Delta_J^{\text{Neu}})^{-1} m = \Pi_J (\Delta^{\text{Neu}})^{-1} m \in \Lambda_J$$

We obtain the **perturbed  $(m, \phi)$ -formulation**, a **nonconforming  $(p, \phi)$ -formulation**:

$$p\text{-}(m, \phi) \quad L \sim \left( \begin{array}{cc} \Delta V \Delta + k^2 (\Delta_J^{\text{Neu}})^{-1} \operatorname{div} V \operatorname{grad} (\Delta_J^{\text{Neu}})^{-1} & k^2 (\Delta_J^{\text{Neu}})^{-1} \operatorname{div} V \operatorname{curl} \\ k^2 \operatorname{rot} V \operatorname{grad} (\Delta_J^{\text{Neu}})^{-1} & k^2 \operatorname{rot} V \operatorname{curl} \end{array} \right)$$

**Analysis as a nonconforming method possible.**



## ( $\phi, m, p, \lambda, \alpha$ ) -formulation

The perturbed p- ( $m, \phi$ ) formulation allows an **equivalent reformulation** as a Galerkin method for a ( $4 \times 4$ ) system that can be considered as a mixed method where the equation ( $p, m$ ) is considered as a constraint and a Lagrange parameter  $\lambda$  is introduced.

To take care of integrability conditions and free constants, one introduces  $\alpha \in \mathbb{C}^3$ .

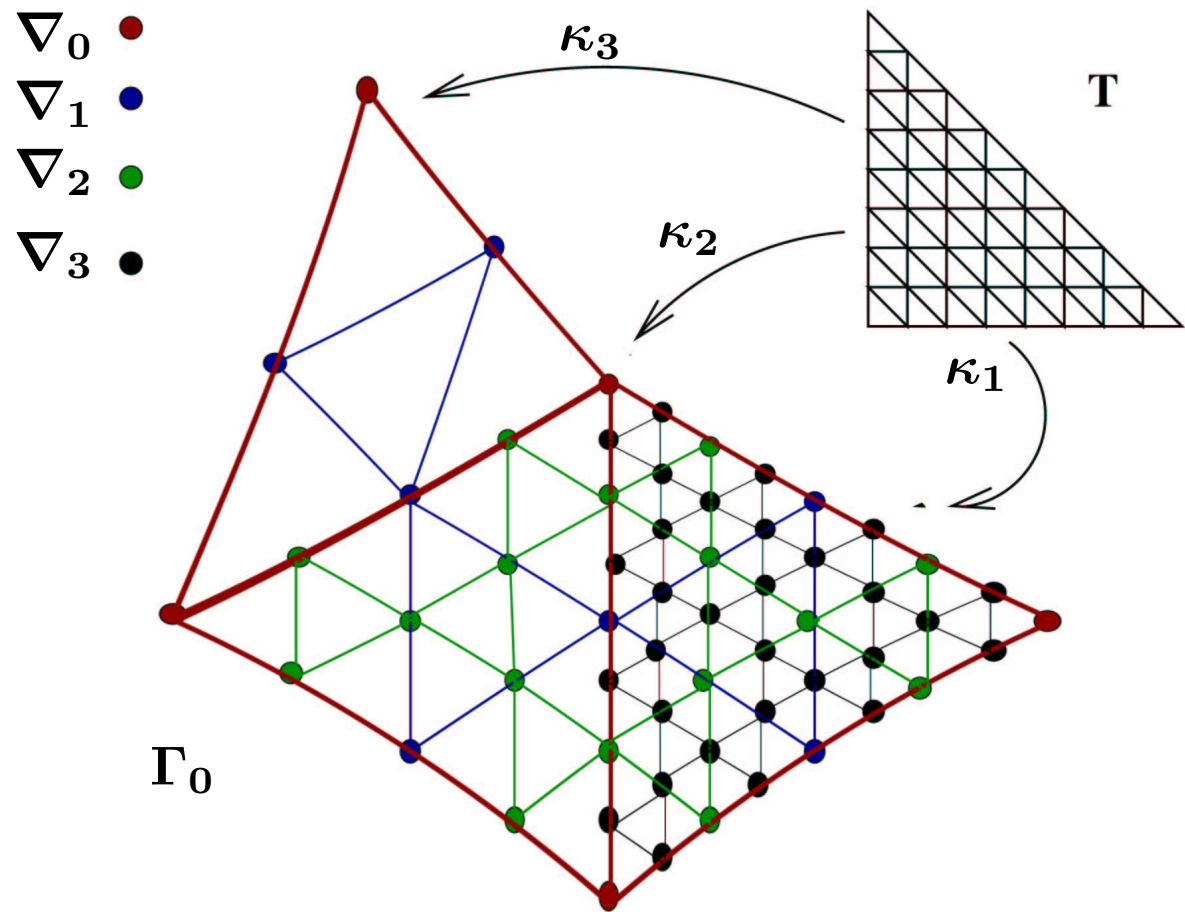
The final result is the system

$$(\phi, m, p, \lambda, \alpha) \quad L \sim \begin{pmatrix} k^2 \text{rot } V \text{ curl} & 0 & k^2 \text{rot } V \text{ grad} & 0 & 0 \\ 0 & V & 0 & 1 & \\ 0 & 1 & -\Delta & 0 & I_3 \\ k^2 \text{div } V \text{ curl} & 0 & k^2 \text{div } V \text{ grad} & -\Delta & \\ 0 & & I_3 & & 0 \end{pmatrix}$$

**Strongly elliptic system. SIOs of order  $\pm 1$ .**

**Energy space:  $\widetilde{H}^{1/2} \times \widetilde{H}^{-1/2} \times H^1 \times H^1 \times \mathbb{C}^3$ . Discretisation by  $\Lambda_J^4 \times \mathbb{C}^3$ .**

# “Nodal” wavelets [Dahmen-Stevenson '99, Rathsfeld]



Triangular patches on  $\Gamma_0$

## “Nodal” wavelets

- Reference triangle  $T$ , refinements  $T^j$ :  $4^j$  triangles, meshsize  $h_j = 2^{-j}$
- Mappings  $\kappa_i : T \rightarrow \Gamma_0, i = 1, \dots, M$ , + continuity conditions  $\rightarrow$   
Coarse triangular patches on  $\Gamma_0$ , nodes  $\Delta_0$
- Refinement level  $j$ :  $\bigcup_{i=1}^M \kappa_i(T^j)$ , nodes  $\Delta_j$
- $P_1$  elements

$$\Lambda_j = \{u \in C^0(\Gamma_0) \mid u \circ \kappa_i|_{T_k^i} \in \mathbb{P}_1; i = 1, \dots, M; k = 1, \dots, 4^j\}$$

- Nodal basis  $\varphi_\tau^j, \tau \in \Delta_j$
- Dual “basis”  $\theta_\tau^{j+1} = \alpha_\tau \varphi_\tau^{j+1} - \beta_\tau \varphi_\tau^j \in \Lambda_{j+1}$ :  $(\varphi_\tau^j, \theta_\sigma^{j+1})_0 = c \delta_{\tau\sigma}$
- Wavelets: difference nodes  $\tau \in \nabla_{j+1} = \Delta_{j+1} \setminus \Delta_j$

$$\psi_\tau = \varphi_\tau^{j+1} - \sum_{\substack{\sigma \\ \text{supp } \varphi_\sigma^j \cap \text{supp } \varphi_\tau^{j+1} \neq \emptyset}} \gamma_{\tau\sigma}^j \theta_\sigma^{j+1}$$

$$\psi_\tau = \varphi_\tau^0, \quad \tau \in \nabla_0 = \Delta_0$$

## Basic properties of the wavelets, closed surface

- **Hierarchy of nodes:**  $\Delta_J = \bigcup_{j=0}^J \nabla_j$
- **Basis of  $\Lambda_J$ :**  $\{\psi_\tau \mid \tau \in \nabla_j, j = 0, \dots, J\}$
- **Support:**  $\tau \in \nabla_j \Rightarrow \text{diam supp } \psi_\tau \sim 2^{-j}$
- **Stability:**  $\forall s \in [-1, 1]$  :

$$\left\| \sum_{j \geq 0} \sum_{\tau \in \nabla_j} c_\tau \psi_\tau \right\|_{H^s(\Gamma)}^2 \sim \sum_{j \geq 0} 2^{2sj} \sum_{\tau \in \nabla_j} |c_\tau|^2$$

- **Approximation property (Jackson) and inverse estimate (Bernstein)**
- **Two vanishing moments:** For  $d = 0, 1$  :

$$|(v, \psi_\tau)_0| \leq C 2^{-(d+2)j} \sup\{|D^\alpha(v \circ \kappa_i)(x)| \mid x \in T; i \leq M; |\alpha| \leq d + 1\}$$

Here  $(u, v)_0 = \sum_i \int_T u \circ \kappa_i \overline{v \circ \kappa_i} dx$  or alternatively,  $(u, v)_0 = \int_\Gamma u \bar{v} ds$

## Estimates for matrix elements

Let  $H$  be a singular integral operator on  $\Gamma$  with kernel  $H(x, y)$  satisfying

$$\forall \alpha, \beta \quad \forall x \neq y \in \Gamma : |D_x^\alpha D_y^\beta H(x, y)| \leq C_{\alpha, \beta} |x - y|^{-(2+r+\alpha+\beta)}$$

Then  $\forall \tau \in \nabla_j, \tau' \in \nabla_{j'}, \text{supp } \psi_\tau \cap \text{supp } \psi_{\tau'} = \emptyset$ :

$$|(H\psi_\tau, \psi_{\tau'})_0| \leq C 2^{-3(j+j')} d_{\tau\tau'}^{-(6+r)}$$

Here  $r = \pm 1$  is the order of the operator and  $d_{\tau\tau'} = \text{dist}(\text{supp } \psi_\tau, \text{supp } \psi_{\tau'})$

The constant  $C$  is independent of the refinement level, but depends on the size of the bounded manifold  $\Gamma$  and on the wave number  $k$ .

This estimate is the basis for a matrix compression scheme based on the distance  $d_{\tau\tau'}$ . One defines a threshold matrix  $(\epsilon_{jj'})_{j, j'=1, \dots, J}$ . For a matrix (block)  $M$  corresponding to the operator  $H$  one defines the compressed matrix  $M^c$  by the rule

$$(\tau \in \nabla_j, \tau' \in \nabla_{j'}) : M_{\tau\tau'}^c = \begin{cases} M_{\tau\tau'} & : d_{\tau\tau'} \leq \epsilon_{jj'} \\ 0 & : d_{\tau\tau'} > \epsilon_{jj'} \end{cases}$$

## Matrix compression, closed surface

Choice of the threshold matrix: Parameters  $K, a, b > 0$ ,

$$\epsilon_{jj'} = K \max\{2^{-j}, 2^{-j'}, 2^{aJ-b(j+j')}\}$$

**Theorem:** Let  $K > 1$ ,  $0 < a < 1$ ,  $b \geq (a + 1)/2$ . Then

$$\#\{M_{\tau\tau'}^c \neq 0\} = \mathcal{O}(J 2^{2J})$$

Define the diagonal order reduction matrix  $D_s = \text{diag}(2^{sj})$

**Theorem:** Let  $M$  be a block corresponding to an operator of order  $r = \pm 1$ . Define  $d = r + 4$ . Let  $K > 1$ ,  $0 < a < 1$ ,  $b = (a + 1)/2$ . Then for all  $s, s' < bd - 2$ , there holds

$$\|D_{-s}(M - M^c)D_{-s'}\|_{\ell^2} \leq C K^{-d} 2^{(r-s-s')J}$$

In our case, we want to choose  $s, s' \in \{0, \pm 1/2\}$  such that  $s + s' \geq r + 1/2$ .

This requires

$$\frac{2}{3} < b = \frac{a + 1}{2} < 1$$

## Wavelets on open surfaces

- Two types of wavelets: **with or without boundary condition.**
- Without boundary condition:  $\Lambda_J, \psi_\tau \dots$  as before
- **With Dirichlet condition** ( $\Gamma_0$  has a polygonal boundary):

**Nodes:**  $\tilde{\Delta}_j = \Delta_j \setminus \partial\Gamma_0, \tilde{\nabla}_j = \tilde{\Delta}_j \setminus \tilde{\Delta}_{j-1} = \nabla_j \setminus \partial\Gamma_0$

**Shape functions:**  $\tilde{\Lambda}_j = \Lambda_j \cap H_0^1(\Gamma_0) = \{\varphi_\tau^j \mid \tau \in \tilde{\Delta}_j\}$

**Wavelets:**  $\tau \in \tilde{\nabla}_j : \tilde{\psi}_\tau \in \tilde{\Lambda}_j \ominus \tilde{\Lambda}_{j-1}$

**Attention:**  $(\Lambda_j \cap H_0^1(\Gamma_0)) \ominus (\Lambda_{j-1} \cap H_0^1(\Gamma_0)) \neq (\Lambda_j \ominus \Lambda_{j-1}) \cap H_0^1(\Gamma_0) !$

**Interior wavelets:**  $\tau \in \tilde{\nabla}_j^0 \Leftrightarrow \forall \sigma \in \partial\Gamma_0 \cap \Delta_j : \text{supp } \tilde{\psi}_\tau \cap \text{supp } \varphi_\sigma^j = \emptyset$

**Boundary layer:**  $\tilde{\nabla}_j^\partial = \tilde{\nabla}_j \setminus \tilde{\nabla}_j^0$ . **Size:**  $\#\tilde{\nabla}_j^\partial \sim 2^j$

**Riesz-Stability:**  $\{\psi_\tau\}$  in  $H^s$  and  $\tilde{H}^{-s}$ ,  $\{\tilde{\psi}_\tau\}$  in  $\tilde{H}^s$  and  $H^{-s}$ ,  $0 \leq s \leq 1$

## Estimates for wavelets on open surfaces

- **Vanishing moments:**

$$|(v, \psi_\tau)| \leq C 2^{-3j} \|v\|_{PW^{2,\infty}(\Gamma_0)} \quad \tau \in \nabla_j$$

$$|(v, \tilde{\psi}_\tau)| \leq C 2^{-3j} \|v\|_{PW^{2,\infty}(\Gamma_0)} + C 2^{-j} \|v\|_{L^\infty(\partial\Gamma_0)} \quad \tau \in \tilde{\nabla}_j$$

- **Matrix elements:**  $\forall \tau \in \tilde{\nabla}_j, \tau' \in \tilde{\nabla}'_j :$

$$\begin{aligned} |(H\psi_\tau, \psi_{\tau'})_0| &\lesssim 2^{-3(j+j')} d_{\tau\tau'}^{-(6+r)} \\ &\quad + 2^{-(j+j')-2\min(j,j')} d_{\tau\tau'}^{-(4+r)} \\ &\quad + 2^{-(j+j')} d_{\tau\tau'}^{-(2+r)} \end{aligned}$$



## Estimates cont'd

- **Compression strategy:**

- All matrix elements where both test and trial functions are close to the boundary are retained
- If both test and trial functions are interior wavelets, one uses the threshold matrix  $\epsilon_{jj'}$  as before
- If one function is interior and the other one in the boundary layer, one uses a second threshold matrix

$$\epsilon_{jj'}^* = K \max\{2^{-j}, 2^{-j'}, 2^{aJ-bj-b^*j'}\}$$

with an additional parameter  $b^*$ .

- **Choice of parameters:**

$$K > 1, \quad \frac{4}{5} < b = \frac{a+1}{2} < 1, \quad b^* = b - 1/2$$

- **Everything then works.**

## Conclusion

One obtains a linear system of  $N \sim 2^{2J}$  equations with

- a simple diagonal preconditioner
- only  $O(N \log N)$  non-zero matrix elements
- quasi-optimal error estimates

$$u = u^1 + u^2, \quad \text{rot } u^1 = 0 = \text{div } u^2$$

$$\|u^1 - u_J^1\|_{H_{\text{div}}^{-1/2}} \lesssim 2^{-J/2}$$

$$\|u^2 - u_J^2\|_0 \lesssim 2^{-J}$$

$$\|\text{div } u - m_J\|_{-1/2} \lesssim 2^{-J/2}$$

Future:

- Implementation
- Comparison with other methods ...